# $\mathbf{P d}_{6} \mathrm{CuB}_{3}$ - a new structure type of borides. $\mathbf{T h}_{7} \mathrm{Fe}_{3}$-type derivative structures in $\mathbf{P d}(\mathbf{P t})-\mathrm{Cu}-\mathrm{B}$ systems. 

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#### Abstract

A missing member of the series of $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type derivative structures, $h$ $\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}$ (unique structure type $\mathrm{Pd}_{6} \mathrm{CuB}_{3}$, space group $P 6_{3} \mathrm{~cm}$, $\mathrm{a}=12.9426(9) \AA$, $\mathrm{c}=4.8697(4) \AA$ ) was obtained from as cast alloys and alloys annealed at $600{ }^{\circ} \mathrm{C}-650{ }^{\circ} \mathrm{C}$. Further substitution of Cu by Pd led to formation of a $\mathrm{Mn}_{7} \mathrm{C}_{3}$-type structure, o$\left(\mathrm{Pd}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathrm{~B}_{3}$ (space group Pnma, $\mathrm{a}=4.8971(2) \AA, \mathrm{b}=7.5353(3) \AA, \mathrm{c}=12.9743(6) \AA$ ). Isotypic $h-\left(\mathrm{Pt}_{0.70} \mathrm{Cu}_{0.30}\right)_{7} \mathrm{~B}_{3}$ was observed in the $\mathrm{Pt}-\mathrm{Cu}-\mathrm{B}$ system as a low temperature phase ( $\mathrm{T} \leq 600{ }^{\circ} \mathrm{C}$ ), whereas the B-filled $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type ( $\mathrm{HT} h$ - $\left(\mathrm{Pt}_{0.717} \mathrm{Cu}_{0.283}\right)_{7} \mathrm{~B}_{3+\mathrm{x}}$, space group $P 6_{3} m c$, $\mathrm{a}=7.4424(12) \AA, \mathrm{c}=4.8549(8) \AA)$ proved to be stable at high temperature. The three structures are built of columns of face connected metal octahedra and columns of metal tetrahedra alternatingly fused by common faces and vertices. Boron atoms are found in trigonal prisms formed by metal atoms; additionally octahedral boron coordination was encountered in HT $h$ $\left(\mathrm{Pt}_{0.717} \mathrm{Cu}_{0.283}\right)_{7} \mathrm{~B}_{3+\mathrm{x}}$. A superconductivity was discovered for $\mathrm{Pt}_{4.9} \mathrm{Cu}_{2.1} \mathrm{~B}_{3}\left(\mathrm{Pd}_{6} \mathrm{CuB}_{3}\right.$-type) and $\mathrm{Pt}_{5.04} \mathrm{Cu}_{1.96} \mathrm{~B}_{3.3}$ (B-filled $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type) below 0.67 and 0.66 K , respectively. Despite the close value of the transition temperature the values of the upper critical field at 0 K differ as 0.37 T and 0.27 T for the two compounds.


## Introduction

The $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type structure (space group $P 6_{3} m c$, $\mathrm{a}=9.85 \AA, c=6.15 \AA$, Th1 in octahedral site $6 c$, Th2 and Th3 in tetrahedral sites $6 c$ and $2 b$ respectively, Fe in trigonal prismatic void in $6 c$ ), ${ }^{[1]}$ which is particularly common among the rare earth - transition metal (M) binary systems, is also adopted by transition metal carbides and borides. Atomic arrangements of the $\mathrm{Th}_{7} \mathrm{Fe}_{3}$ type carbides, so-called polytypes of the Eckstrom-Adcock phase, ${ }^{[2]}$ attract continuous research interest from the materials science community because of the significance of these compounds in manufacturing of chromium steels, high-alloyed cast irons and other engineering materials. ${ }^{[3-6]}$ The $\mathrm{M}_{7} \mathrm{C}_{3}$ carbide structures usually exhibit many planar defects; ${ }^{[7]}$ an extensive investigation of these defects in $\mathrm{M}_{7} \mathrm{C}_{3}(\mathrm{M}=\mathrm{Cr}, \mathrm{Fe})$ crystals by electron diffraction and resulting structures were presented by Kowalski (1985). ${ }^{[8]}$ Nevertheless, the hexagonal $P 6_{3} m c^{[9]}$ and orthorhombic Pnma ${ }^{[10]}$ unit cells $\left(\mathrm{Th}_{7} \mathrm{Fe}_{3^{-}}\right.$and $\mathrm{Mn}_{7} \mathrm{C}_{3}$-type structure, respectively) have been commonly considered for $\mathrm{M}_{7} \mathrm{C}_{3}$ transition metal carbides. Very recently, a new orthorhombic $o-\mathrm{Fe}_{7} \mathrm{C}_{3}{ }^{[11]}$ phase ( $\mathrm{a}=4.5202 \AA, \mathrm{~b}=13.7572 \AA, \mathrm{c}=11.974 \AA$ ) isotypic with the $\mathrm{Ca}_{7} \mathrm{Au}_{3}$-type (space group $\left.P b c a\right)^{[12]}$ has been derived from high-pressure single crystal X-ray diffraction data. In contrast to transition metal carbide systems, only few binary transition metal borides were found to crystallize in the $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type structure: $\mathrm{Ru}_{7} \mathrm{~B}_{3},{ }^{[9]}$ $\mathrm{Rh}_{7} \mathrm{~B}_{3},{ }^{[13]} \mathrm{Re}_{7} \mathrm{~B}_{3},{ }^{[13]}$ and $\mathrm{Tc}_{7} \mathrm{~B}_{3}{ }^{[14]}$. Besides, a disordered variant of $\mathrm{Th}_{7} \mathrm{Fe}_{3}$ has been reported
recently for a metastable $\mathrm{Ni}_{7} \mathrm{~B}_{3}$ compound. ${ }^{[15]} \mathrm{M} / \mathrm{M}^{\prime}$ substitution has been studied for $\mathrm{Re}_{7} \mathrm{~B}_{3}$ by Kuz'ma et al. $\left(\mathrm{M}^{\prime}=\mathrm{Co}, \mathrm{Ni}, \mathrm{Nb}, \mathrm{Hf}, \mathrm{Ta}\right)^{[16]}$ and for $\mathrm{Ru}_{7} \mathrm{~B}_{3}$ and $\mathrm{Rh}_{7} \mathrm{~B}_{3}$ by Fokwa et al. $\left(\mathrm{M}^{\prime}=\right.$ $\mathrm{Cr}, \mathrm{Mn}, \mathrm{Fe}, \mathrm{Co}, \mathrm{Ni})^{[17]}$. Comparative studies of the structures of $\mathrm{Rh}_{7} \mathrm{~B}_{3}, \mathrm{Ru}_{7} \mathrm{~B}_{3}$ and $\mathrm{Re}_{7} \mathrm{~B}_{3}$ binaries in relation to $\mathrm{M}_{7} \mathrm{C}_{3}$ carbides ${ }^{[18]}$ elucidated the hexagonal symmetry for $\mathrm{M}_{7} \mathrm{~B}_{3}$ and revealed the absence of either superstucture reflections or diffuse intensities thus confirming the structure of borides as suggested by Aronsson ${ }^{[9,13]}$.

Considering (i) the fact that boron and carbon tend to form homologues compounds at low non-metal contents (e.g. $\mathrm{Cr}_{23} \mathrm{C}_{6}$-type, $\mathrm{Fe}_{3} \mathrm{C}$-type, etc.) ${ }^{[19]}$ as well as (ii) the ambiguity in solving complex boride structures based on powder diffraction data, the absence of $\mathrm{Th}_{7} \mathrm{Fe}_{3}-$ type derivative structures among borides was intriguing. Our recent work into $\mathrm{Pd}(\mathrm{Pt})-\mathrm{Cu}-\mathrm{B}$ systems ${ }^{[20,21]}$ exemplified the potential of boride and carbide systems to exhibit common structural arrangements, which still have to be synthesized. Herein we report the discovery of a novel series of ternary metal borides of the $\mathrm{Th}_{7} \mathrm{Fe}_{3}$ family, among which the hexagonal $h$ $\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}$ adopts a new structure type, the $\mathrm{Pd}_{6} \mathrm{CuB}_{3}$ type, whereas orthorhombic $o$ $\left(\mathrm{Pd}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathrm{~B}_{3}\left(\mathrm{Mn}_{7} \mathrm{C}_{3}\right.$-type) has never been observed before among transition metal borides. The hexagonal structure discovered is apparently so far a missing member of the structural series exhibiting a volume ratio of $1\left(\mathrm{Th}_{7} \mathrm{Fe}_{3}\right.$-type $): 2\left(\mathrm{Mn}_{7} \mathrm{C}_{3}\right.$-type $): 3\left(\mathrm{Pd}_{6} \mathrm{CuB}_{3}-\right.$ type) : 4 ( $\mathrm{Ca}_{7} \mathrm{Au}_{3}$-type). We also report in the current work on a B-filled $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type, which was observed for the first time from as cast alloys in the $\mathrm{Pt}-\mathrm{Cu}-\mathrm{B}$ system while the samples annealed at $600^{\circ} \mathrm{C}$ showed isotypism with the novel $\mathrm{Pd}_{6} \mathrm{CuB}_{3}$-type. The structures of all three new borides were determined by single-crystal X-ray diffraction and confirmed from Rietveld refinements of powder X-ray diffraction data; we should note that the current study on the orthorhombic $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-derivative boride is the first one performed from single crystal intensity data since the discovery of this structural arrangement from Weissenberg photographs in $1969^{10}$. The relationships between the new structure type and the three derivatives of $\mathrm{Th}_{7} \mathrm{Fe}_{3}{ }^{-}$ type are discussed. We present herein as well the electrical transport properties down to 0.3 K for four compounds including the evaluation of upper critical fields of the new superconductors $\mathrm{Pt}_{4.9} \mathrm{Cu}_{2.1} \mathrm{~B}_{3}\left(\mathrm{Pd}_{6} \mathrm{CuB}_{3}\right.$-type) and $\mathrm{Pt}_{5.04} \mathrm{Cu}_{1.96} \mathrm{~B}_{3.3}$ (B-filled $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type); the results obtained have provided impetus for further investigation of physical properties of these compounds which will be reported in a forthcoming paper. Because the structural studies of transition metal carbides are challenging owing to experimental difficulties arising from the highly faulted nature of $\mathrm{MM}_{7}{ }_{7} \mathrm{C}_{3}$ crystals ${ }^{[7,8]}$, the obtained crystallographic data on homologues boride structures are significant for further development of the important technological materials.

## Results and Discussion

Crystal structure of HT $\boldsymbol{h}-\left(\mathbf{P t}_{0.717} \mathbf{C u}_{0.283}\right)_{7} \mathbf{B}_{3+\mathrm{x}}, \mathbf{x}=\mathbf{0 . 2 8}$. The positions of heavy atoms and B1 obtained from structure refinement in the space group $P 6_{3} m c$ from X-ray single crystal data (Tables 1-2) are consistent with the $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type ${ }^{[1]}$. Accordingly, metal atoms in the structure of platinum copper boride (Figure 1) are distributed within the three thorium atom sites, two of which randomly accommodate Pt and Cu atoms ( M 2 in $6 c$ and M 3 in $2 b$, respectively). The Th1 octahedral site is solely occupied by platinum ( Pt 1 in $6 c$ ); about $30 \%$ of $\left[\mathrm{Pt} 1_{6}\right]$ octahedra are centered by additional boron atoms ( B 2 in $2 a$ ). Coordination polyhedra of atoms are shown in Figure1 Supporting Information. B1 is found at the iron atom site inside a trigonal prism formed by metal atoms. The B-metal atom distances in trigonal prisms are ranging within $2.00 \AA$ and $2.19 \AA$ (Table 1 Supporting Information); two additional M2 atoms are located against the rectangular faces of the prism at a distance of 0.268 nm . The shortest contact of $2.00 \AA$ between boron and Pt 1 in the $\left[\mathrm{B}_{2} \mathrm{Pt}_{6}\right]$ octahedron is not unusual considering partial occupancy of the boron atom site. Atoms M2 and M3 are coordinated by 15 atoms. The coordination number of Pt 1 has increased to 16 as compared to the corresponding metal atom in boride representatives of $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type, e.g. Ru in $\mathrm{Ru}_{7} \mathrm{~B}_{3}$ due to two extra B 2 atoms delivered by a new boron site centering the $\left[\mathrm{Pt}_{6}\right]$ octahedra.


Figure 1. Crystal structure of $\mathrm{HT} h-\left(\mathrm{Pt}_{0.717} \mathrm{Cu}_{0.283}\right)_{7} \mathrm{~B}_{3+\mathrm{x}}, \mathrm{x}=0.28$ as arrangement of B filled trigonal prisms (grey; in online version - orange), B filled octahedra (light grey; in online version - yellow) and empty tetrahedra (dark grey; in online version - green).

${ }^{[a]}$ atom coordinates are standardized using the program Structure Tidy ${ }^{[22]}$

| Atom | Site | x | y | Z | Occupancy | $\mathrm{U}_{\mathrm{eq}}{ }^{[a, b]}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| HT $\boldsymbol{h}$-( $\left.\mathbf{P t}_{0.717} \mathrm{Cu}_{0.283}\right)_{7} \mathbf{B}_{3+\mathrm{x}}, \mathbf{x}=\mathbf{0 . 2 8}$ |  |  |  |  |  |  |
| Pt1 | 6 c | 0.12645(6) | 0.87355 | 0.237(9) | 1.00 | 0.0126(3) |
| M2 | 6 c | 0.5477(1) | 0.4523(1) | 0.069(9) | $0.446(3) \mathrm{Pt} 2+0.554(3) \mathrm{Cu} 2$ | 0.0164(4) |
| M3 | $2 b$ | 1/3 | 2/3 | 0.048(9) | $0.681(4) \mathrm{Pt} 3+0.319(4) \mathrm{Cu} 3$ | 0.0165(5) |
| B1 | $6 c$ | 0.813(2) | 0.187(2) | 0.342(11) | 1.00 | 0.014(5) |
| B2 | $2 a$ | 0 | 0 | $0.000^{\text {c }}$ | 0.28(7) | 0.014(5) |
| $\boldsymbol{h}-\left(\mathbf{P d}_{0.864} \mathbf{C u}_{0.136}\right)_{7} \mathbf{B}_{3}$ |  |  |  |  |  |  |
| Pd1 | 12d | 0.20828(8) | 0.54320(8) | 0.3674(2) | 1.00 | 0.0071(2) |
| M2 | $12 d$ | 0.2075(1) | 0.3315(2) | 0.1765(3) | $0.627(4) \mathrm{Pd} 2+0.373(4) \mathrm{Cu} 2$ | 0.0095(3) |
| Pd3 | $6 c$ | 0.3352(2) | 0 | 0.1901(5) | 1.00 | 0.0088(2) |
| Pd4 | $6 c$ | 0.1233(1) | 0 | $0.0000^{\text {c }}$ | 1.00 | 0.0073(3) |
| M5 | 6 c | 0.5498(1) | 0 | 0.2186(4) | $0.793(4) \mathrm{Pd} 5+0.207(4) \mathrm{Cu} 5$ | 0.0088(4) |
| B1 | $12 d$ | 0.324(1) | 0.474(1) | 0.470(3) | 1.00 | 0.011(1) |
| B2 | 6 c | 0.193(2) | 0 | 0.399(4) | 1.00 | 0.011(1) |
| $o-\left(\mathrm{Pd}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathrm{~B}_{3}$ |  |  |  |  |  |  |
| Pd1 | 8d | 0.06598(6) | 0.56369(4) | 0.31094(3) | 1.00 | 0.0070(1) |
| M2 | $8 d$ | 0.24897(7) | 0.06486(5) | 0.01827(3) | $0.786(3) \mathrm{Pd} 2+0.214(3) \mathrm{Cu} 2$ | 0.0090(1) |
| Pd3 | 4 c | 0.23498(9) | 1/4 | 0.41479(4) | 1.00 | 0.0067(1) |
| Pd4 | 4 c | 0.42506(9) | $1 / 4$ | 0.62647(3) | 1.00 | 0.0067(1) |
| M5 | $4 c$ | 0.2081(1) | 1/4 | 0.20098(4) | $0.932(3) \mathrm{Pd} 5+0.068(3) \mathrm{Cu} 5$ | 0.0073(1) |
| B1 | $8 d$ | 0.461(1) | 0.0317(7) | 0.3510(4) | 1.00 | 0.011(1) |
| B2 | 4 c | 0.033(1) | 1/4 | 0.5609(6) | 1.00 | 0.014(1) |


| Table 1 Supporting Information. Selected interatomic distances [ $\AA$ ] |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
| $\mathbf{H T ~} \boldsymbol{h}$-( $\left(\mathrm{Pt}_{0.717} \mathrm{Cu}_{0.283}\right)_{7} \mathbf{B}_{3+\mathrm{x}}, \mathbf{x}=\mathbf{0 . 2 8}{ }^{a}$ |  |  |  |  |
| Pt1- B2 2.00(3) | M2- 2B1 2.19(5) | M3-3B1 2.13(4) | B1- Pt1 2.07(3) | B2-3 Pt1 2.00(3) |
| - B2 2.07(3) | - 2M2 2.655(3) | -3M2 2.766(3) | - M3 2.13(4) | - 3 Pt1 2.07(3) |
| - B1 2.07(3) | -2B1 2.68(3) | -3M2 2.777(2) | - 2Pt1 2.19(3) |  |
| - 2B1 2.19(3) | - 2M2 2.722(1) | -3Pt1 2.82(3) | - 2M2 2.19(5) |  |
| - M3 2.82(3) | - M3 2.766(3) | -3M2 2.97(6) | - 2M2 2.68(3) |  |
| - 2Pt1 2.823(2) | - M3 2.777(2) |  |  |  |
| - 2M2 2.83(4) | - 2Pt1 2.83(4) |  |  |  |
| - 2M2 2.84(2) | - 2Pt1 2.84(2) |  |  |  |
| - 4Pt1 2.92(6) | - M3 2.97(6) |  |  |  |
| $h-\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathbf{B}_{3}{ }^{\text {b }}$ |  |  |  |  |
| Pd1- B1 2.11(1) | M2- B2 2.17(2) | Pd3- B2 2.10(2) | Pd4- B2 2.14(2) | M5- 2B1 2.19(2) |
| - B1 2.16(1) | - B1 2.22(1) | -2B1 2.15(1) | - 2B2 2.25(2) | - 2M2 2.729(2) |
| - B1 2.26(2) | - B1 2.51(2) | -2M2 2.711(4) | - 2Pd4 2.765(2) | - Pd3 2.763(3) |
| -2Pd1 2.786(2) | - B2 2.60(2) | - M5 2.736(3) | - M2 2.824(2) | - 2M5 2.755(3) |
| - M2 2.824(2) | - 2M2 2.664(2) | - M5 2.780(3) | - M2 2.824(2) | - Pd3 2.780(3) |
| - Pd3 2.829(3) | - Pd3 2.711(4) | -2Pd1 2.829(3) | - 2M2 2.827(2) | - 2Pd1 2.834(1) |
| - M5 2.834(1) | - M5 2.729(2) | -2M2 2.875(2) | - Pd3 2.895(3) | - 2Pd1 2.900(2) |
| - M2 2.889(2) | - M2 2.780(3) | - Pd4 2.895(3) | - 4Pd4 2.911(1) | - Pd3 2.973(3) |
| - M5 2.900(2) | - Pd4 2.824(2) | - M5 2.973(3) |  |  |
| -2Pd1 2.901(1) | - Pd1 2.824(2) | - 2M2 2.985(3) | B2-Pd3 2.10(2) | B1-Pd1 2.11(1) |
| -2Pd1 2.935(2) | - Pd4 2.827(2) |  | - Pd4 2.14(2) | - Pd3 2.15(1) |
|  | - Pd3 2.875(2) |  | - 2M2 2.17(2) | - Pd1 2.16(1) |
|  | - Pd1 2.889(2) |  | - 2Pd4 2.25(2) | - M5 2.19(2) |
|  | - Pd3 2.985(3) |  | - 2M2 2.60(2) | - M2 2.22(1) |
|  |  |  |  | - Pd1 2.26(2) |
|  |  |  |  | - M2 2.51(2) |
| $\left.\boldsymbol{o - ( ~} \mathbf{P d}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathbf{B}_{3}{ }^{\boldsymbol{c}}$ |  |  |  |  |
| Pd1- B1 2.127(5) | M2- B1 2.221(5) | Pd3- B2 2.139(8) | Pd4- B2 2.101(7) | M5- 2B1 2.151(5) |
| - B2 2.229(6) | - B2 2.221(6) | -2B1 2.148(5) | - 2B1 2.215(5) | - 2M2 2.729(2) |
| - B1 2.280(5) | - B1 2.509(5) | -2M2 2.7272(5) | - 2Pd1 2.7751(5) | - 2M5 2.7592(7) |
| - Pd4 2.7751(5) | - B2 2.660(3) | - M5 2.7610(7) | - 2Pd2 2.8265(5) | - Pd3 2.7610(7) |
| - Pd1 2.8078(4) | - M2 2.6695(5) | - M5 2.7772(7) | - 2Pd2 2.8855(5) | - Pd3 2.7772(7) |
| - Pd2 2.8385(6) | - M2 2.6880(5) | -2Pd1 2.8439(4) | - Pd3 2.8999(7) | - 2Pd1 2.8472(5) |
| - Pd3 2.8439(4) | - Pd3 2.7272(5) | -2M2 2.8923(5) | - 2Pd1 $2.9005(5)$ | - 2Pd1 2.9468(4) |
| - Pd5 2.8472(5) | - M5 2.7579(6) | - Pd4 2.8999(7) | - 2Pd1 2.9737(5) | - Pd3 2.9857(7) |
| - Pd2 2.8738(5) | - M2 2.7902(5) | - M5 2.9857(7) |  |  |
| - Pd4 2.9005(5) | - Pd4 2.8265(5) | - 2M2 3.0060(6) | B2-Pd4 2.101(7) | B1-Pd1 2.127(5) |
| -2Pd1 2.9148(5) | - Pd1 2.8385(6) |  | - Pd3 2.139(8) | - Pd3 2.148(5) |
| - Pd5 2.9468(4) | - Pd1 2.8738(5) |  | - 2M2 2.221(6) | - M5 2.151(5) |
| - Pd4 2.9737(5) | - Pd4 2.8855(5) |  | - 2Pd1 2.229(6) | - Pd4 2.215(5) |
|  | - Pd3 2.8923(5) |  | - 2M2 2.660(3) | - M2 2.221(5) |
|  | - Pd3 3.0060(6) |  |  | - Pd1 2.280(5) |
|  |  |  |  | - M2 2.509(5) |

${ }^{a} \mathrm{M} 2=0.446(3) \mathrm{Pt} 2+0.554(3) \mathrm{Cu} 2, \mathrm{M} 3=0.681(4) \mathrm{Pt} 3+0.319(4) \mathrm{Cu} 3$
${ }^{b} \mathrm{M} 2=0.627$ (4) Pd2 $+0.373(4) \mathrm{Cu} 2, \mathrm{M} 5=0.793(4) \mathrm{Pd} 5+0.207$ (4) Cu 5
${ }^{c} \mathrm{M} 2=0.786(3) \mathrm{Pd} 2+0.214(3) \mathrm{Cu} 2, \mathrm{M} 5=0.932(3) \mathrm{Pd} 5+0.068(3) \mathrm{Cu} 5$
Crystal structure of $\boldsymbol{h}-\left(\mathbf{P d}_{0.864} \mathbf{C u}_{0.136}\right)_{7} \mathbf{B}_{3}$ - a new $\mathbf{P d}_{6} \mathbf{C u B}_{3}$-type. The compound represents a unique type of the $\mathrm{Th}_{7} \mathrm{Fe}_{3}$ derivative structures where palladium atoms occupy two octahedral sites ( Pd 1 in $12 d$ and Pd 4 in $6 c$ ) and one tetrahedral site ( Pd 3 in $6 c$ ) while the remaining two tetrahedral sites are randomly populated by Pd and Cu (M2 in $12 d$ and M 5 in $6 c$ ) (Tables 1-2, Figure 2). Trigonal prisms formed by metal atoms with one (and two) additional M2 atoms against their rectangular faces accommodate B1 and B2 respectively
(Table 1 Supporting Information, Figure 1k,1 Supporting Information). The $\left[\operatorname{Pd} 1_{6}\right]$ and $\left[\operatorname{Pd} 4_{6}\right]$ octahedra are not occupied by boron in the current crystal. Both Pd1 and Pd4 are coordinated by 11 metal atoms and 3 borons. Similarly to M3 in $\mathrm{HT} h-\left(\mathrm{Pt}_{0.717} \mathrm{Cu}_{0.283}\right)_{7} \mathrm{~B}_{3+\mathrm{x}}$ of B -filled $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type, the Pd3 have 15 neighbors; M2 and M5 are bounded to 15 and 13 atoms respectively.


Figure 2. Crystal structure of $h-\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}$ as arrangement of B -filled trigonal prisms (grey; in online version - orange), empty octahedra (light grey; in online version - yellow) and empty tetrahedra (dark grey; in online version - green).

| Atom | Site | $\mathrm{U}_{11}$ | $\mathrm{U}_{22}$ | $\mathrm{U}_{33}$ | $\mathrm{U}_{23}$ | $\mathrm{U}_{13}$ | $\mathrm{U}_{12}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| HT $\boldsymbol{h}$-( $\left(\mathrm{P}_{0.717} \mathrm{Cu}_{0.283}\right)_{7} \mathbf{B}_{3+\mathrm{x}}, \mathrm{x}=0.28$ |  |  |  |  |  |  |  |
| Pt1 | 6 c | 0.0095(4) | 0.0095(4) | 0.0204(5) | -0.0008(2) | 0.0008(2) | 0.0059(3) |
| M2 | 6 c | 0.0132(6) | 0.0132(6) | 0.0216(9) | -0.0042(3) | $0.0042(3)$ | 0.0056(6) |
| M3 | $2 b$ | 0.00942 (6) | 0.00942 (6) | 0.0307(12) | 0 | 0 | 0.0047(3) |
| $h$ - $\left(\mathbf{P d}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathbf{B}_{3}$ |  |  |  |  |  |  |  |
| Pd1 | $12 d$ | $0.0066(4)$ | 0.0068(4) | 0.0071(4) | 0 | -0.0002(5) | 0.0026(4) |
| M2 | $12 d$ | 0.0084(7) | 0.0107(6) | 0.0094(6) | 0.0000(3) | 0.0006(5) | 0.0045(6) |
| Pd3 | 6 c | 0.0080(5) | 0.0082(9) | 0.0102(4) | 0 | 0.0005(3) | 0.0041(5) |
| Pd4 | $6 c$ | 0.0068(5) | 0.0080(7) | 0.0075(7) | 0 | -0.0008(6) | 0.0040(3) |
| M5 | 6 c | 0.0090(6) | 0.0096(9) | 0.0079(8) | 0 | $0.0017(5)$ | 0.0048(5) |
| ${ }_{\boldsymbol{o}}-\left(\mathbf{P d}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathbf{B}_{3}$ |  |  |  |  |  |  |  |
| Pd1 | $8 d$ | 0.0076(2) | 0.0070(2) | 0.0063(2) | 0.0009(1) | -0.0004(1) | 0.0006(1) |
| M2 | $8 d$ | 0.0096(2) | 0.0088(2) | 0.0085(2) | -0.0010(1) | 0.0002(1) | 0.0001(1) |
| Pd3 | $4{ }_{c}$ | 0.0075(2) | 0.0069(2) | 0.0058(2) | 0 | $0.0007(2)$ | 0 |
| Pd4 | $4{ }^{\text {c }}$ | 0.0082(2) | 0.0069(2) | 0.0051(2) | 0 | -0.0002(2) | 0 |
| M5 | 4 c | 0.0057(2) | 0.0104(2) | 0.0057(2) | 0 | 0.0006(2) | 0 |

Crystal structure of $\boldsymbol{o}-\left(\mathbf{P d}_{0.93} \mathbf{C u}_{0.07}\right)_{7} \mathbf{B}_{\mathbf{3}}$. Figure 3 shows the crystal structure of a new boride of $\mathrm{Mn}_{7} \mathrm{C}_{3}$-type ${ }^{[10]}$ (Tables 1-2). As viewed along the shortest axis, the structure is composed of two columns of empty $\left[\mathrm{Pd}_{4} \mathrm{Pd}_{2}\right]$ octahedra and four columns of edge- and face-connected empty [ $\mathrm{M} 2{ }_{2} \mathrm{M} 5 \mathrm{Pd} 3$ ] tetrahedra (Figure 3). Coordination polyhedra of atoms and interatomic distances are given in Figure 1, Supporting Information and Table 1, Supporting Information. Alike boron atoms in the $h-\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}$ structure, B 1 and B 2 are found in trigonal prisms formed by metal atoms. Coordination numbers of Pd 1 and Pd 4 and shapes of their coordination polyhedra correspond to that of Pd 1 and Pd 4 in the $h-\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}$ structure as well as resembles the coordination polyhedron of Pt 1 in the $\mathrm{HT} h-\left(\mathrm{Pt}_{0.717} \mathrm{Cu}_{0.283}\right)_{7} \mathrm{~B}_{3+\mathrm{x}}$ structure, if not taking into account the surplus boron atoms delivered by populated centers of octahedra of the latter one. Analogously to Pd 3 in the $h-\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}$ structure of
 and comprises the displacement of the B 2 atom as compared to M 3 in HT $h$ $\left(\mathrm{Pt}_{0.717} \mathrm{Cu}_{0.283}\right)_{7} \mathrm{~B}_{3+\mathrm{x}}$. The shape of the coordination sphere of $\mathrm{M} 5\left(\left[\mathrm{M}_{5} \mathrm{~B}_{2} \mathrm{M}_{4} \mathrm{Pd}_{7}\right]\right.$ is analogous to that of M5 in the $h-\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}$ structure.


Figure 3. Crystal structure of $o-\left(\mathrm{Pd}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathrm{~B}_{3}$ as arrangement of B -filled trigonal prisms (grey; in online version - orange), empty octahedra (light grey; in online version - yellow) and empty tetrahedra (dark grey; in online version - green).

A closer examination of coordination spheres of M 2 in $\mathrm{HT} h-\left(\mathrm{Pt}_{0.717} \mathrm{Cu}_{0.283}\right)_{7} \mathrm{~B}_{3+\mathrm{x}}, h-$ $\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}$ and $o-\left(\mathrm{Pd}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathrm{~B}_{3}$ shows that, despite of their identical coordination numbers $(\mathrm{CN}=15)$, the arrangements of surrounding atoms are different. Whilst in the B-filled $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type structure, the four boron atoms coordinating the M 2 site exhibit a planar orientation (Figure 4a), in the two remaining structures the boron atoms are arranged
tetrahedrally with respect to the central ones and realize, together with neighboring atoms, different geometries of the $\left[\mathrm{M} 2 \mathrm{Pd} 1_{2} \mathrm{Pd} 4_{2} \mathrm{M} 2_{3} \mathrm{M} 5_{1} \mathrm{Pd}_{3}{ }_{3} \mathrm{~B} 1_{2} \mathrm{~B} 2_{2}\right]$ polyhedra in the $h$ $\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}$ and $o-\left(\mathrm{Pd}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathrm{~B}_{3}$ structures (Figures 4b, c, respectively).


Figure 4. M2 coordination polyhedra in $\mathrm{HT} h-\left(\mathrm{Pt}_{0.717} \mathrm{Cu}_{0.283}\right)_{7} \mathrm{~B}_{3+\mathrm{x}}$ (a), $h-\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}$ (b) and $o-\left(\mathrm{Pd}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathrm{~B}_{3}$ (c) emphasizing the disposition of boron atoms in the polyhedral units. Atoms codes correspond to those given in Figures 1-3.


Figure 1 Supporting Information. Coordination polyhedra of atoms in the structures HT $h$ $\left(\mathrm{Pt}_{0.717} \mathrm{Cu}_{0.283}\right)_{7} \mathrm{~B}_{3+\mathrm{x}}\left(\mathrm{a}-\mathrm{Pt} 1, \mathrm{~b}-\mathrm{M} 2, \mathrm{c}-\mathrm{M} 3\right.$, d-B1, e-B2), $h-\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}(\mathrm{f}-\mathrm{Pd} 1, \mathrm{~h}-\mathrm{M} 2, \mathrm{j}-$ Pd3, g-Pd4, i-M5, k-B1, l-B2) and $o-\left(\mathrm{Pd}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathrm{~B}_{3}(\mathrm{~m}-\mathrm{Pd} 1$, o-M2, q-Pd3, n-Pd4, p-M5, rB1, s-B2).

Structure relationships. All three new boride structures are closely related with the $\mathrm{Th}_{7} \mathrm{Fe}_{3}{ }^{-}$ type ${ }^{[1]}$. According to common description, ${ }^{[7,23]}$ the $\mathrm{Th}_{7} \mathrm{Fe}_{3}$ structure is composed of infinite columns of face-connected empty [ $\mathrm{Th} 1_{6}$ ] octahedra and columns of empty [ $\mathrm{Th} 2_{3} \mathrm{Th} 3_{1}$ ] tetrahedra both running infinitely along the short axis (Figure 5a-d). Tetrahedra condense alternatingly via common vertices and faces. Iron atoms are located between the columns in trigonal prisms formed by adjacent "free" faces of octahedra and tetrahedra. Every second tetrahedron in the column participates in the formation of the trigonal prisms (Figure 5d) and only half of the "free" faces of octahedra can form the bases of trigonal prisms (Figure 5c). Consequently, all prisms are sloping in one direction with respect to the short axis. Triads of [ $\mathrm{FeTh} 1_{3} \mathrm{Th} 2_{2} \mathrm{Th} 3_{1}$ ] trigonal prisms, associated with successive octahedra along a column differ by $60^{\circ}$ in orientation (Figure 5a). The trigonal prisms share edges and vertices to form a three-dimensional framework. In the parent representative of the $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type boride family, $\mathrm{Ru}_{7} \mathrm{~B}_{3},{ }^{[9]} \mathrm{Ru}$ substitutes for Th whereas B takes the atom site of Fe .

The HT $h$ - $\left(\mathrm{Pt}_{0.717} \mathrm{Cu}_{0.283}\right)_{7} \mathrm{~B}_{3+\mathrm{x}}$ structure ( B -filled variant of the $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type) and $o$ $\left(\mathrm{Pd}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathrm{~B}_{3}\left(\mathrm{Mn}_{7} \mathrm{C}_{3}\right.$-type $)$ differ in the relative positioning of the columns of octahedra one of which is displaced and rotated by $60^{\circ}$ around the short axis (slab B, Figure 6b) with respect to the other (slab A, Figure 6b) in the $o-\left(\mathrm{Pd}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathrm{~B}_{3}$ structure, while in the structure of $\mathrm{HT} h-\left(\mathrm{Pt}_{0.717} \mathrm{Cu}_{0.283}\right)_{7} \mathrm{~B}_{3+\mathrm{y}}$ the columns of octahedra are identical and located at the same height along $z$ (slab A, Figure 6a). Besides the $\mathrm{Th}_{7} \mathrm{Fe}_{3}-$ and $\mathrm{Mn}_{7} \mathrm{C}_{3}$-type structures, a hexagonal unit cell exhibiting identical columns of octahedra but manifesting the doubled $a$ and $b$ lattice parameters (suggested space groups $P 6_{3} m c$ or $P 31 c$, $\mathrm{a}=1.387 \mathrm{~nm}, \mathrm{c}=0.463 \mathrm{~nm}$ ) was hitherto also suggested ${ }^{[24]}$ to describe the $\mathrm{M}_{7} \mathrm{C}_{3}$-carbides, however this structure proposal was superseded by further studies. ${ }^{[10,25,26]}$
$h$ - $\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}\left(\mathrm{Pd}_{6} \mathrm{CuB}_{3}\right.$-type) reported in the current study, exhibits a new structural arrangement which maintains the stoichiometry and structural units of the $\mathrm{Th}_{7} \mathrm{Fe}_{3}{ }^{-}$ type structure. The unit cell of the new structure is composed of a total of three columns of face-connected metal octahedra running infinitely along the short axis, two of which are formed by $\mathrm{Pd} 1(12 d)$ and one is built by $\mathrm{Pd} 4(6 c)$. Thus, whilst in the $o-\left(\mathrm{Pd}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathrm{~B}_{3}$ unit cell every second column of the existing two is moved around and along the shorter axis, $h$ $\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}$ exhibited the arrangement of two displaced columns (slab B-B formed by Pd1) vs one unmoved (slab A formed by Pd4) (Figure 6c).


Figure 5. Boron filled trigonal metal prisms (grey; in online version - yellow) and their clustering around the columns of empty metal octahedra (left open) (a, c) and tetrahedra (dark grey; in online version - brown) (b, d) in the structure of $\mathrm{Th}_{7} \mathrm{Fe}_{3}$ extending infinitively along the short axis. Arrangements of B1 and B2 filled trigonal metal prisms (light grey and grey respectively; in online version - yellow and pink respectively) around a column of empty tetrahedra (brown) in the $h$ - $\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}\left(\mathrm{Pd}_{6} \mathrm{CuB}_{3}\right.$-type) structure (e). Atoms codes correspond to those given in Figures 1-3.

In the structure of $\mathrm{HT} h$ - $\left(\mathrm{Pt}_{0.717} \mathrm{Cu}_{0.283}\right)_{7} \mathrm{~B}_{3+\mathrm{x}}\left(\mathrm{B}\right.$-filled $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type) every subsequent tetrahedron in the column offers its three "free" faces to form the bases of boron filled trigonal prisms (Figure 5d), while in the remaining structures described herein, every tetrahedron in the column participates in the formation of boron filled trigonal prisms (Figure 5e) alternatingly contributing either one or two faces. Thus the trigonal prisms attached to the tetrahedral columns are alternatively sloping up and down in the $o-\left(\mathrm{Pd}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathrm{~B}_{3}\left(\mathrm{Mn}_{7} \mathrm{C}_{3}-\right.$ type) and $h$ - $\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}\left(\mathrm{Pd}_{6} \mathrm{CuB}_{3}\right.$-type $)$ structures.

The empty octahedron in the $o-\left(\mathrm{Pd}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathrm{~B}_{3}$ structure $\left(\mathrm{Mn}_{7} \mathrm{C}_{3}\right.$-type) is built of two kinds of Pd atoms and the resulting octahedral columns are surrounded by both B1- and B2filled trigonal prisms. The composition and local environment of octahedra in the $h$ $\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}$ structure $\left(\mathrm{Pd}_{6} \mathrm{CuB}_{3}\right.$-type $)$ are different from $o-\left(\mathrm{Pd}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathrm{~B}_{3}$. The columns are composed of either [ $\left.\mathrm{Pd} 1_{6}\right]$ or $\left[\mathrm{Pd} 4_{6}\right]$ octahedra; each of them is surrounded by a different kind of borons. Thus the trigonal prisms around B 1 and B 2 in the $o-\left(\mathrm{Pd}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathrm{~B}_{3}$ structure form a complex three-dimensional framework being connected through common edges and vertices (Figure 7b), whereas in the $h$ - $\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}$ structure the B 1 -filled trigonal prisms share edges and single corners to form a three-dimensional framework which envelops the columns of $\left[P d 4_{6}\right]$ octahedra and associated triades of $\left[\mathrm{B} 2 \mathrm{Pd} 4_{3} \mathrm{M} 2_{2} \mathrm{Pd} 3_{1}\right]$ trigonal prisms (Figure 7a).


Figure 6. Arrangement of octahedra and adjacent atoms in $\mathrm{HT} h-\left(\mathrm{Pt}_{0.717} \mathrm{Cu}_{0.283}\right)_{7} \mathrm{~B}_{3+\mathrm{x}}(\mathrm{a}), h-$ $\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}$ (b), $o-\left(\mathrm{Pd}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathrm{~B}_{3}$ (c) and $\mathrm{Ca}_{7} \mathrm{Au}_{3}$ (d) structures, respectively. For better comparison, hexagonal structures are presented in orthohexagonal settings. Metal atoms which do not form octahedra are omitted. Atoms codes correspond to those given in Figure 13. For $\mathrm{Ca}_{7} \mathrm{Au}_{3}(\mathrm{~d})-\mathrm{Ca}$ is dark grey and Au is grey (in online version dark blue and red respectively).


Figure 7. The frameworks of boron filled trigonal prisms (B1-grey, B2 - dark grey; in online version yellow and pink respectively) in the structures $h-\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}\left(\mathrm{Pd}_{6} \mathrm{CuB}_{3}\right.$-type, a) and $o-\left(\mathrm{Pd}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathrm{~B}_{3}\left(\mathrm{Mn}_{7} \mathrm{C}_{3}\right.$-type, b).

All structures reported here, ( $\mathrm{HT} h$ - $\left(\mathrm{Pt}_{0.717} \mathrm{Cu}_{0.283}\right)_{7} \mathrm{~B}_{3+\mathrm{x}}$ ( B -filled $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type), o$\left(\mathrm{Pd}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathrm{~B}_{3}\left(\mathrm{Mn}_{7} \mathrm{C}_{3}\right.$-type $), h-\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}\left(\mathrm{Pd}_{6} \mathrm{CuB}_{3}\right.$-type $)$ as well as the $\mathrm{Ca}_{7} \mathrm{Au}_{3}-$ type ${ }^{[12]}$ structure which exhibits two A fragments vs two B (Figure 6d), maintain the octahedra-to-tetrahedra ratio equal to 2 and their unit cell volumes relate as 1:2:3:4, respectively, considering the existence of $\mathrm{Ca}_{7} \mathrm{Au}_{3}$-type related platinum copper boride described herein below exhibiting the lattice parameters $a=4.9248 \AA, b=14.563 \AA, c=12.920$ A.

Formation of compounds and site preferences. The investigations of compounds with the $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type derivative structures have been focused hitherto on chromium iron carbides alloyed with additional transition metals due to dramatic influence of the eutectic $\mathrm{MM}^{\prime}{ }_{7} \mathrm{C}_{3}$ carbides on the microstructure and mechanical properties of high-chromium white cast irons. ${ }^{[27]}$ According to structural studies, the orthorhombic $\mathrm{M}_{7} \mathrm{C}_{3}$ phase of $\mathrm{Mn}_{7} \mathrm{C}_{3}$-type was found for synthesis temperatures above $1200{ }^{\circ} \mathrm{C}$. At lower temperatures the diffraction patterns indicated a hexagonal structure (space group $P 6_{3} m c, \mathrm{Th}_{7} \mathrm{Fe}_{3}$-type). ${ }^{[28]}$ A complete solid solution has been reported between $\mathrm{Cr}_{7} \mathrm{C}_{3}$ and $\mathrm{Mn}_{7} \mathrm{C}_{3}{ }^{[29]}$ as well as between $\mathrm{Cr}_{7} \mathrm{C}_{3}$ and $\mathrm{Fe}_{7} \mathrm{C}_{3},{ }^{[30]}$ solubility of Ni and W in $\mathrm{Cr}_{7} \mathrm{C}_{3}$ was found to be smaller: $\mathrm{Cr}_{7-\mathrm{x}} \mathrm{Ni}_{\mathrm{x}} \mathrm{C}_{3}, \mathrm{x}=0-0.5{ }^{[29]}$ and $\mathrm{Cr}_{7-\mathrm{x}} \mathrm{W}_{\mathrm{x}} \mathrm{C}_{3}, \mathrm{x}=0.4^{[31]}$. Despite of high industrial interest, there is a lack of reliable data on solubility and preferential positions for atoms of alloying elements in the multicomponent $\left(\mathrm{M}, \mathrm{M}^{\prime}\right)_{7} \mathrm{C}_{3}$ carbide systems $\left(\mathrm{M}, \mathrm{M}^{\prime}=\mathrm{V}, \mathrm{Cr}, \mathrm{Mn}, \mathrm{Fe}, \mathrm{Co}, \mathrm{W}\right) .{ }^{[27,32]}$ In the $\mathrm{Fe}-\mathrm{Cr}-\mathrm{V}-\mathrm{C}$ type alloys, vanadium substitutes either for Fe or Cr depending on the $\mathrm{Fe} / \mathrm{Cr}$ ratio. ${ }^{[33]}$ Recent investigations on the Co-Cr-C sub-system showed that cobalt stabilizes the orthorhombic $\mathrm{Mn}_{7} \mathrm{C}_{3}$-type structure preferring to occupy the M2 and M5 sites in the orthorhombic cell, ${ }^{[34]}$ however the theoretical studies based on ab initio cohesive energies and lattice inversion method suggested the hexagonal $P 6_{3} / m c$ structure as the most stable. ${ }^{[35]}$ A comparison between the obtained energies furthermore showed that at low content, Co atoms prefer to occupy the chromium sites at the vertices of tetrahedra (2b). After increasing the amount of $\mathrm{Co}\left(\mathrm{Cr}_{7-\mathrm{x}} \mathrm{Co}_{\mathrm{x}} \mathrm{C}_{3}, \mathrm{x} \geq 1\right)$, the additional cobalt atoms substitute chromium in tetrahedral $6 c$ and octahedral $6 c$ sites; moreover, the calculations indicated that despite the energy needed for cobalt atoms to occupy octahedral sites is lower, the structure with this atom site substitution pattern (Co atoms in tetrahedral $2 b$ and octahedral $6 c$ ) become unstable if $\mathrm{x}>1.35$. Thereby at higher Co content, the octahedral site is unaffected, whereas two tetrahedral sites accommodate the $\mathrm{Cr} / \mathrm{Co}$ mixtures. ${ }^{[35]}$

In the $\mathrm{Pd}-\mathrm{Cu}-\mathrm{B}$ system, both high Pd content $h$ - $\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}$ and $o-$ $\left(\mathrm{Pd}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathrm{~B}_{3}$ phases form directly from the melt and follow the substitution pattern reported for $\mathrm{Co}_{6.60} \mathrm{Cr}_{0.51} \mathrm{C}_{2.89}{ }^{[34]}$ where the only atom sites forming the adjacent faces of tetrahedra in the columns are prone to $\mathrm{M} / \mathrm{M}^{\prime}$ mixing. As inherited from X-ray powder diffraction studies and EDX, the $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type derivative palladium copper borides span over a rather wide concentration range within $\mathrm{Pd}_{58.5} \mathrm{Cu}_{11.5} \mathrm{~B}_{30}$ to $\mathrm{Pd}_{68.5} \mathrm{Cu}_{1.5} \mathrm{~B}_{30}$ at $650{ }^{\circ} \mathrm{C}$. A hexagonal structure with a tripled unit cell of $\mathrm{Pd}_{6} \mathrm{CuB}_{3}$-type exists from $\mathrm{Pd}_{58.5} \mathrm{Cu}_{1.5} \mathrm{~B}_{30}$ to $\mathrm{Pd}_{63.5} \mathrm{Cu}_{6.5} \mathrm{~B}_{30} \quad(\mathrm{a}=12.9412(1)-12.9640(2) \quad \AA, \quad \mathrm{c}=4.8697(1)-4.8830(1) \quad \AA, \quad \mathrm{V}=706.29(1)-$ $710.72(2) \AA^{3}$ ); increase of Pd amount stabilizes the orthorhombic $\mathrm{Mn}_{7} \mathrm{C}_{3}$-type structure $\left(\mathrm{Pd}_{65} \mathrm{Cu}_{5} \mathrm{~B}_{30}-\mathrm{Pd}_{68.5} \mathrm{Cu}_{1.5} \mathrm{~B}_{30}, \quad \mathrm{a}=4.8884(2)-4.9116(1) \quad \AA, \quad \mathrm{b}=7.5272(2)-7.5585(1) \quad \AA\right.$, $\mathrm{c}=12.9697(2)-12.9733(1) \AA, \mathrm{V}=477.23(2)-481.625(2) \AA^{3}$ (Figure 2 Supporting Information). As delivered by single crystal and powder X-ray diffraction data refinements, the Wyckoff sites concerned with $\mathrm{Pd} / \mathrm{Cu}$ substitution in the Cu -rich side of the $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type derivative structure series are tetrahedral $12 d$ and $6 c$ of the tripled hexagonal unit cell yielding 55.3\% $\mathrm{Pd} 2+44.7 \% \mathrm{Cu} 2$ and $73.7 \% \mathrm{Pd} 5+26.3 \% \mathrm{Cu} 5$ levels of occupancy, respectively (Rietveld refinement of X-ray powder diffraction data). The Pd-richest terminus amounts appr. 1.5 at. \% Cu and exhibits the only one tetrahedral site affected by $\mathrm{Pd} / \mathrm{Cu}$ substitution (92.1\% $\mathrm{Pd} 2+7.9 \%$ Cu 2 in $8 d$ of orthorhombic $\mathrm{Mn}_{7} \mathrm{C}_{3}$-type unit cell) (Rietveld refinement of X-ray powder diffraction data). Although the structures of the Pd-poor and Pd-rich phases are closely related, there is no continuous remodeling of the structures upon variation in Pd content in consistency with the fact that no direct crystallographic group-subgroup relation exists. The critical amount of Cu for destabilization of large hexagonal structure and formation of the orthorhombic structure is within 6.5 and $5 \mathrm{at} . \%$ according to X-ray diffraction data and EDX analysis of two-phase alloys (Figure 3 Supporting Information).

According to EDX measured $\mathrm{Pt} / \mathrm{Cu}$ ratios and X -ray data analysis, the $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type derivative structures in $\mathrm{Pt}-\mathrm{Cu}-\mathrm{B}$ as cast alloys were found to cover the range $\mathrm{Pt}_{0.72}$ $\left.{ }_{\mathrm{y}} \mathrm{Cu}_{0.28+\mathrm{y}}\right)_{7} \mathrm{~B}_{3+\mathrm{x}}(\mathrm{y} \cong 0-0.16, \mathrm{x} \cong 0-0.3)$. The Pt-rich end (HT $h-\left(\mathrm{Pt}_{0.72-\mathrm{y}} \mathrm{Cu}_{0.28+\mathrm{y}}\right)_{7} \mathrm{~B}_{3+\mathrm{x}}, \mathrm{y} \cong 0-0.07$ ) forms a hexagonal structure of B-filled $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type which on annealing at $600{ }^{\circ} \mathrm{C}$ decomposes to the low temperature form of $\mathrm{Pd}_{6} \mathrm{CuB}_{3}$-type (Figure 4 Supporting Information, Table 3 Supporting Information). Strong irregularities in variations of lattice parameters obtained from Rietveld refinement of X-ray powder data for HT $h$ - $\left(\mathrm{Pt}_{0.72-\mathrm{y}} \mathrm{Cu}_{0.28+\mathrm{y}}\right)_{7} \mathrm{~B}_{3+\mathrm{x}}$ within Pt-rich and Cu -rich boundaries depending on boron content with respect to EDX $\mathrm{Pt} / \mathrm{Cu}$ ratios support our results obtained from single crystal X-ray diffraction data on surplus boron
amount in the structure as a consequence of incorporation of boron atoms into the octahedral voids of the $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type unit cell.


Figure 2 Supporting Information. Rietveld refinement of X-ray powder diffraction intensity data of $\mathrm{Pd}_{60} \mathrm{Cu}_{10} \mathrm{~B}_{30}\left(\mathrm{Pd}_{6} \mathrm{CuB}_{3}\right.$-type structure) and $\mathrm{Pd}_{68} \mathrm{Cu}_{2} \mathrm{~B}_{30}\left(\mathrm{Mn}_{7} \mathrm{C}_{3}\right.$-type structure $)$ annealed at $650{ }^{\circ} \mathrm{C}$.


Figure 3 Supporting Information. Rietveld X-ray powder diffraction refinement of $\mathrm{Pd}_{64} \mathrm{Cu}_{6} \mathrm{~B}_{30}$ alloy annealed at $650{ }^{\circ} \mathrm{C}$ revealing the existence of orthorthombic $\mathrm{Mn}_{7} \mathrm{C}_{3}$-type and hexagonal $\mathrm{Pd}_{6} \mathrm{CuB}_{3}$-type phases.


Figure 4 Supporting Information. Rietveld refinement of X-ray powder diffraction intensity data of as cast $\mathrm{Pt}_{50} \mathrm{Cu}_{18} \mathrm{~B}_{32}$ sample ( B -filled $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type, x value fixed in correspondence with results of SC XRD) and $\mathrm{Pt}_{49} \mathrm{Cu}_{21} \mathrm{~B}_{30}$ sample annealed at $600{ }^{\circ} \mathrm{C}\left(\mathrm{Pd}_{6} \mathrm{CuB}_{3}\right.$-type $)$
$\mathrm{Cu} / \mathrm{Pt}$ mixtures in the structures of both HT- and LT-phases reside predominantly at all tetrahedral sites while octahedral sites exhibit only minor replacement of platinum by copper (Table 3 Supporting Information). A similar tendency in site occupancies has been observed for the $\mathrm{Rh}_{7-\mathrm{x}} \mathrm{M}_{\mathrm{x}} \mathrm{B}_{3}(\mathrm{M}=\mathrm{Cr}, \mathrm{Mn}, \mathrm{Fe}, \mathrm{Co}, \mathrm{Ni})^{[17]}$ phases of $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type where M was found to mix with rhodium at only one tetrahedral $6 c$ site at $\mathrm{x} \leq 1$ but after increasing the amount of M , particularly of iron in $\mathrm{FeRh}_{6} \mathrm{~B}_{3}$ to produce, for example, the $\mathrm{Fe}_{1.3} \mathrm{Rh}_{5.7} \mathrm{~B}_{3}$, the additional iron atoms substitute rhodium in the tetrahedral $2 b$ site, thereby leaving only one octahedral $6 c$ site unaffected. For the copper-rich end of $\left(\mathrm{Pt}_{0.72-\mathrm{y}} \mathrm{Cu}_{0.28+\mathrm{y}}\right)_{7} \mathrm{~B}_{3+\mathrm{x}}$ range $(\mathrm{y} \cong 0.07-0.16)$, the similarities in the X-ray powder patterns with the $\mathrm{Ca}_{7} \mathrm{Au}_{3}$-type ${ }^{[12]}$ structure have been elucidated, e.g. $a=4.9248(4) \AA, b=14.563(1) \AA, c=12.920(1) \AA$ in space group $P b c a$ for the $\mathrm{Pt}_{44} \mathrm{Cu}_{26} \mathrm{~B}_{30}$ alloy annealed at $600{ }^{\circ} \mathrm{C}$, however the poor values of reliability factors of Rietveld refinement ( $\mathrm{R}_{\mathrm{F}}=0.10, \mathrm{R}_{\mathrm{I}}=0.16$ ) did not provide sufficient proofs for the correct choice of the structure model.

Our verification of the boundary Pd-B binary system in the relevant concentration and temperature ranges showed good agreement with literature data revealing the existence of three phases, namely $\mathrm{Pd}_{2} \mathrm{~B}, \mathrm{Pd}_{5} \mathrm{~B}_{2}$ and $\mathrm{Pd}_{3} \mathrm{~B}$ as well as the absence of $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type structure ${ }^{[36,37]}$. Traces of an unknown phase noticed on prolonged annealing at $600{ }^{\circ} \mathrm{C}(480$ hours), has obviously indicated the formation of a low temperature unknown ${ } \mathrm{Pd}_{5} \mathrm{~B}$ " structure. The recent studies on the Pt-B binary boundary ${ }^{[38]}$ yielded no formation of a $\mathrm{Th}_{7} \mathrm{Fe}_{3}{ }^{-}$ type phase as well. Since no homologues binary palladium and platinum borides have been found to exist, the formation of the ternary phases is apparently a consequence of an electronic stabilizing effect of alloying element atoms.

| Table 3 Supporting Information. Results of X-ray powder Reetveldrefinement of $\mathrm{LT} h$ - $\left.\mathrm{Pt}_{0.70} \mathrm{Cu}_{0.30}\right)_{7} \mathrm{~B}_{3}$ |  |
| :---: | :---: |
| structure type | $\mathrm{Pd}_{6} \mathrm{CuB}_{3}$ |
| composition (Rietved refinement) | $\mathrm{Pt}_{4.9} \mathrm{Cu}_{2.1} \mathrm{~B}_{3}$ |
| composition (EPMA) | $\mathrm{Pt}_{4.85} \mathrm{Cu}_{2.15} \mathrm{~B}_{3}$ |
| space group | $\mathrm{Pb}_{3} \mathrm{~cm}$, no. 185 |
| lattice parameters [ $\AA$ ] | $a=12.8881$ (1), $c=4.9579(1)$ |
| volume [ $\AA^{3}$ ] | 713.19(1) |
| $\rho_{\text {calc }}\left[\mathrm{g} \mathrm{cm}^{-3}\right], Z$ | 15.60, 6 |
| range for data collection | $8^{\circ}<\theta<100^{\circ}$ |
| reflections in refinement | 190 |
| number of variables | 61 |
| reliability factors ${ }^{b}$ | $\mathrm{R}_{\mathrm{F}}=0.020, \mathrm{R}_{\mathrm{I}}=0.023, \mathrm{R}_{\text {exp }}=0.016, \chi^{2}=3.08$ |
| M1 in 12d ${ }^{\text {a }}$; | $x=0.2061(1) ; \mathrm{y}=0.5402(1) ; \mathrm{z}=0.3811(2)$; |
| occ.; $\mathrm{Biso}^{\text {is }}{ }^{\text {c }}$ | 0.93(1)Pt1+0.07(1)Cu1;0.68(4) |
| M2 in 12d; | $x=0.2071(1) ; \mathrm{y}=0.3333(2) ; \mathrm{z}=0.1762(3)$; |
| occ.; $\mathrm{B}_{\text {iso }}$ | 0.46 (1) Pt2+0.54(1)Cu2; 0.72(3) |
| M3 in 6c; | $x=0.3376(2) ; \mathrm{z}=0.2150(4)$; |
| occ.; $\mathrm{B}_{\text {iso }}$ | 0.68(1)Pt3+32(1)Cu3; 0.60(4) |
| M4 in 6 c; | $x=0.1239(1) ; \mathrm{z}=0.0000^{d}$; |
| occ.; $\mathrm{B}_{\text {iso }}$ | $0.85(1) \mathrm{Pt} 4+0.15(1) \mathrm{Cu} 4 ; 0.46$ (2) |
| M5 in 6 c; | $x=0.5580(2) ; \mathrm{z}=0.2608(4)$; |
| occ.; $\mathrm{B}_{\text {iso }}$ | 0.55(1)Pt5+0.45(1)Cu5; 0.48(3) |
| B1 in 12d; | $x=0.324(1) ; \mathrm{y}=0.474(1) ; \mathrm{z}=0.470(3)$; |
| occ.; $\mathrm{B}_{\text {iso }}$ | $1.00 \mathrm{B1} ; 0.8{ }^{\text {d }}$ |
| B2 in 6 c; | $x=0.193$ (2); $\mathrm{z}=0.399(4)$; |
| occ.; $\mathrm{B}_{\text {iso }}$ | $1.00 \mathrm{~B} 2 ; 0.8{ }^{\text {d }}$ |

${ }^{a}$ atom coordinates are standardized using the program Structure Tidy ${ }^{22},{ }^{b} \mathrm{R}_{\mathrm{F}}=\Sigma\left|\mathrm{F}_{\mathrm{o}}-\mathrm{F}_{\mathrm{c}}\right| / \Sigma \mathrm{F}_{\mathrm{o}}, \mathrm{R}_{\mathrm{I}}=\Sigma\left|\mathrm{I}_{0}-\mathrm{I}_{\mathrm{c}}\right| / \Sigma \mathrm{I}_{0}$, $\left.R_{\text {exp }}=\left[(N-P+C) / \Sigma w_{i y} y^{2}{ }_{\mathrm{o}}\right)\right]^{1 / 2}, \chi^{2}=\left(R_{\mathrm{wp}} / R_{e}\right)^{2},{ }^{c}$ isotropic atomic displacement parameters $\left[\AA^{2}\right],{ }^{d}$ fixed parameters

Electrical Resistivity. Concluding from resistivity data, the compounds $\mathrm{Pt}_{4.9} \mathrm{Cu}_{2.1} \mathrm{~B}_{3}$ ( $\mathrm{Pd}_{6} \mathrm{CuB}_{3}$-type), $\mathrm{Pt}_{5.04} \mathrm{Cu}_{1.96} \mathrm{~B}_{3.3}$ (B-filled $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type), $\mathrm{Pd}_{6.504} \mathrm{Cu}_{0.496} \mathrm{~B}_{3}\left(\mathrm{Mn}_{7} \mathrm{C}_{3}\right.$-type) and $\mathrm{Pd}_{6.047} \mathrm{Cu}_{0.953} \mathrm{~B}_{3}\left(\mathrm{Pd}_{6} \mathrm{CuB}_{3}\right.$-type) are metallic in the temperature range from 1 to 300 K (Figure 8). Below 1 K , however $\mathrm{Pt}_{4.9} \mathrm{Cu}_{2.1} \mathrm{~B}_{3}$ and $\mathrm{Pt}_{5.04} \mathrm{Cu}_{1.96} \mathrm{~B}_{3.3}$ exhibit a superconducting transition at 0.67 and 0.66 K , respectively. $\mathrm{Pd}_{6.504} \mathrm{Cu}_{0.496} \mathrm{~B}_{3}$ and $\mathrm{Pd}_{6.047} \mathrm{Cu}_{0.953} \mathrm{~B}_{3}$, on the other hand, show superconductivity at temperatures close to the experimental limit.

The normal state resistivity curves for $\mathrm{Pd}_{6.504} \mathrm{Cu}_{0.496} \mathrm{~B}_{3}$ and $\mathrm{Pd}_{6.047} \mathrm{Cu}_{0.953} \mathrm{~B}_{3}$ were analyzed in scope of the Bloch-Grüneisen-Mott relation,

$$
\begin{equation*}
\rho_{B-G-M}=\rho_{0}+C \frac{T^{5}}{\theta_{D}{ }^{6}} \int_{0}^{\theta_{D} / T} \frac{x^{5}}{\left(e^{x}-1\right)\left(1-e^{-x}\right)} d x-A_{M} T^{3} \tag{1}
\end{equation*}
$$

where $\rho_{0}$ is a residual resistivity at $0 \mathrm{~K}, \theta_{D}$ is a Debye temperature and a Mott-Jones term $\left(-A_{M} T^{3}\right)^{[39]}$ is added to the conventional Bloch -Grüneisen relation to account for the corrections due to scattering of conduction electrons on a narrow $d$ band in the vicinity of the Fermi energy. Debye temperatures $\theta_{D}$, residual resistivity values $\rho_{0}$ and Mott coefficients $A_{M}$ obtained as the results of the fit of Eqn. 1 to the appropriate experimental data are presented in Table 3, together with the residual resistivity ratios $\left(\operatorname{RRR}=\rho_{300} / \rho_{N}\right.$, where $\rho_{N}$ is the value of the resistivity just above the superconducting transition) for the corresponding compounds.


Figure 8. The electrical resistivity (a) and $\rho / \rho_{300 K}$ (b) of $\mathrm{Pt}_{4.9} \mathrm{Cu}_{2.1} \mathrm{~B}_{3},\left(\mathrm{Pd}_{6} \mathrm{CuB}_{3}\right.$-type) (1), $\mathrm{Pt}_{5.04} \mathrm{Cu}_{1.96} \mathrm{~B}_{3.3} \quad\left(\mathrm{~B}\right.$-filled $\quad \mathrm{Th}_{7} \mathrm{Fe}_{3}$-type) (2), $\mathrm{Pd}_{6.504} \mathrm{Cu}_{0.496} \mathrm{~B}_{3} \quad\left(\mathrm{Mn}_{7} \mathrm{C}_{3}\right.$-type) (3) and $\mathrm{Pd}_{6.047} \mathrm{Cu}_{0.953} \mathrm{~B}_{3}\left(\mathrm{Pd}_{6} \mathrm{CuB}_{3}\right.$-type) (4). Solid lines correspond to models described in text.

Electrical resistivity of both $\mathrm{Pt}_{4.9} \mathrm{Cu}_{2.1} \mathrm{~B}_{3}$ and $\mathrm{Pt}_{5.04} \mathrm{Cu}_{1.96} \mathrm{~B}_{3.3}$, complimentary to the low RRR, demonstrates stronger curvature in the temperature region from 30 to 100 K and a tendency to saturate at higher temperatures. This feature no longer allows a description in terms of Eqn.1. Instead a good agreement was found between the experimental data of both compounds for temperatures from 10 to 300 K and a phenomenological model proposed by Woodard and Cody ${ }^{[40]}$ to account for the resistivity of A15 superconductor $\mathrm{Nb}_{3} \mathrm{Sn}$,

$$
\rho=\rho_{0}+A T+B e^{(-C / T)}
$$

where $\rho_{0}$ is the resistivity at low temperatures, $A$ characterizes the growth of the resistivity at high temperatures and the exponential term is included to describe the unexpected rapid temperature dependence of the scattering probability in the low temperature region. The parameters obtained as a result of the least-square fits are listed in Table 3.

|  | $\rho_{0},[\mu \Omega \mathrm{~cm}]$ | $\theta_{D}[\mathrm{~K}]$ | $A_{M}\left[10^{-8} \mu \Omega \mathrm{~cm} / \mathrm{K}^{3}\right]$ | $A[\mu \Omega \mathrm{~cm} / \mathrm{K}]$ | $C[\mathrm{~K}]$ | RRR |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{Pt}_{4.9} \mathrm{Cu}_{2.1} \mathrm{~B}_{3}$ | 196.2 | - | - | 0.025 | 109.8 | 1.04 |
| $\mathrm{Pt}_{5.04} \mathrm{Cu}_{1.96} \mathrm{~B}_{3.3}$ | 152.6 | - | - | 0.0081 | 74.0293 | 1.03 |
| $\mathrm{Pd}_{6.504} \mathrm{Cu}_{0.496} \mathrm{~B}_{3}$ | 40.9 | 216.6 | 1.6 | - | - | 1.89 |
| $\mathrm{Pd}_{6.077} \mathrm{Cu}_{0.953} \mathrm{~B}_{3}$ | 121.6 | 203.9 | 3.7 | - | - | 1.30 |

To get a deeper insight into the behavior of $\mathrm{Pt}_{4.9} \mathrm{Cu}_{2.1} \mathrm{~B}_{3}$, and $\mathrm{Pt}_{5.04} \mathrm{Cu}_{1.96} \mathrm{~B}_{3.3}$ their upper critical fields were studied and analyzed. The curves of the electrical resistivity obtained at various externally applied magnetic fields are displayed in Figure 9. An application of magnetic fields reduces the superconducting transition temperature, which becomes zero at the upper critical magnetic field $\mu_{0} H_{C 2}(0)$. Whereas for $\mathrm{Pt}_{4.9} \mathrm{Cu}_{2.1} \mathrm{~B}_{3}$, the superconducting transition can still be observed at magnetic fields as high as 0.27 T , the 0.25 T resistivity curve of $\mathrm{Pt}_{5.04} \mathrm{Cu}_{1.96} \mathrm{~B}_{3.3}$ reveals only the weakly noticeable onset of the transition, while the transition itself happens below the lowest temperature reachable in our setup. For both compounds the superconducting transition is completely suppressed at magnetic fields higher then 0.5 T .

The temperature dependent upper critical magnetic field was analyzed using a Werthamer, Helfand, and Hohenberg (WHH) model ${ }^{[41]}$ that considers both orbit-depairing and Pauli mechanisms for the destruction of the Cooper pairs. Corresponding to this model, the temperature dependence of the upper critical field is determined by the transition temperature $T_{C}$, the slope of the upper critical field at the transition temperature $\left.\frac{d\left[\mu_{0} H_{C 2}(T)\right]}{d T}\right|_{T c}$, the Maki parameter $\alpha_{M}$ and spin-orbit coupling $\lambda_{s o}$, where $\left.\frac{d\left[\mu_{0} H_{C 2}(T)\right]}{d T}\right|_{T_{c}}$ can be directly measured from a linear approximation of the experimental data near the transition temperature and the Maki parameter can be calculated as $\alpha_{M}=\frac{3 e^{2} \mathrm{~h} \gamma \rho_{0}}{2 m \pi^{2} k_{B}^{2}}$. In our case the lack of heat capacity data makes it impossible to directly calculate $\alpha_{M}$, rather it was obtained as
$\alpha_{M}=5.3 \times\left. 10^{-5} \frac{-d\left[\mu_{0} H_{C 2}(T)\right]}{d T}\right|_{T_{C}}{ }^{[42]}$ (see Table 4). It is important to note that the influence of $\lambda_{s o}$ on the temperature dependence of the upper critical field is negligible if $\alpha_{M}$ is small. The value of $\lambda_{s o}$ in our study is about 10 for both compounds; the values of the upper critical fields at $\mu_{0} H_{C 2}(0)$ are 0.37 and 0.27 T for $\mathrm{Pt}_{4.9} \mathrm{Cu}_{2.1} \mathrm{~B}_{3}$ and $\mathrm{Pt}_{5.04} \mathrm{Cu}_{1.96} \mathrm{~B}_{3.3}$, respectively.
$\mathrm{Pd}_{6.504} \mathrm{Cu}_{0.496} \mathrm{~B}_{3}$ and $\mathrm{Pd}_{6.047} \mathrm{Cu}_{0.953} \mathrm{~B}_{3}$ in the low temperature limit hint to the onset of superconductivity by a small resistivity drop (Figure 9d).

| Table 4. Parameters describing the behavior of the temperature dependent upper critical magnetic <br> field for $\mathrm{Pt}_{4.9} \mathrm{Cu}_{2.1} \mathrm{~B}_{3}$ and $\mathrm{Pt}_{5.04} \mathrm{Cu}_{1.96} \mathrm{~B}_{3.3}$ |  |  |  |
| :--- | :--- | :--- | :--- |
| $T_{C}[\mathrm{~K}]$ | $\left.\frac{d\left[\mu_{0} H_{C 2}(T)\right]}{d T}\right\|_{T c}\left[\mu_{0}(T / K)\right]$ | $\alpha_{M}$ | $\lambda_{s o}$ |
|  | 0.67 | -0.8 |  |
| $\mathrm{Pt}_{4.9} \mathrm{Cu}_{2.1} \mathrm{~B}_{3}$ | -0.6 | 0.42 | $\sim 10$ |
| $\mathrm{Pt}_{5.04} \mathrm{Cu}_{1.96} \mathrm{~B}_{3.3}$ | 0.66 |  | 0.32 |



Figure 9. Evolution of the electrical resistivity under the influence of magnetic fields of $\mathrm{Pt}_{4.9} \mathrm{Cu}_{2.1} \mathrm{~B}_{3}$ (a), $\mathrm{Pt}_{5.04} \mathrm{Cu}_{1.96} \mathrm{~B}_{3.3}$ (b) and their temperature dependent upper critical fields $\mu_{0} H_{C 2}(T)(\mathrm{c})$. Electrical resistivity of $\mathrm{Pd}_{6.504} \mathrm{Cu}_{0.496} \mathrm{~B}_{3}$ and $\mathrm{Pd}_{6.047} \mathrm{Cu}_{0.953} \mathrm{~B}_{3}$ in the low temperature region. The drop in $\rho(T)$ is attributed to the onset of superconductivity (d).

## Conclusions

The current study expanded the knowledge on the $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type derivative structures introducing a new member of the boride family which exhibits a hitherto unknown structure type, $\mathrm{Pd}_{6} \mathrm{CuB}_{3}$ (space group $P 6_{3} c m, \quad \mathrm{a}=1.29426(9) \mathrm{nm}, \quad \mathrm{c}=0.48697(4) \mathrm{nm}$ for $h$ $\left.\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}\right)$. Besides, at slightly higher palladium content a first boride representative of the orthorhombic $\mathrm{Mn}_{7} \mathrm{C}_{3}$-type has been found to exist in the $\mathrm{Pd}-\mathrm{Cu}-\mathrm{B}$ system: o$\left(\mathrm{Pd}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathrm{~B}_{3}$, (space group Pnma, $\mathrm{a}=0.48971$ (2) $\mathrm{nm}, \mathrm{b}=0.75353$ (3) $\mathrm{nm}, \mathrm{c}=1.29743(6) \mathrm{nm}$ ). Furthermore, in the $\mathrm{Pt}-\mathrm{Cu}-\mathrm{B}$ system, single crystal data of $\mathrm{HT} h-\left(\mathrm{Pt}_{0.717} \mathrm{Cu}_{0.283}\right)_{7} \mathrm{~B}_{3+\mathrm{x}}$ (space group $P 6_{3} m c, a=0.74424(12) \mathrm{nm}, \mathrm{c}=0.48549(8) \mathrm{nm}$ ) have proven a new B-filled structure variant of $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type which exhibited a partial B-occupancy of octahedral [ $\mathrm{Pt}_{6}$ ] voids whereas the low-temperature phase was observed to form the $\mathrm{Pd}_{6} \mathrm{CuB}_{3}$-type structure (space group $P 6_{3} c m$, $\mathrm{a}=1.28996$ (1) nm, $\mathrm{b}=0.49635(1) \mathrm{nm}$ for LT $\left.h-\left(\mathrm{Pt}_{0.70} \mathrm{Cu}_{0.30}\right)_{7} \mathrm{~B}_{3}\right)$ from Rietveld refinement of X-ray powder data of alloys annealed at $600^{\circ} \mathrm{C}$. The core part of the all three structures is built of columns of face-connected metal octahedra and columns of condensed alternatively via common vertices and faces metal tetrahedra. Boron atoms fill the trigonal prismatic voids within the metal atom frameworks in the three structures as well as partially fill the octahedral voids in the structure of $\mathrm{HT} h$ - $\left(\mathrm{Pt}_{0.717} \mathrm{Cu}_{0.283}\right)_{7} \mathrm{~B}_{3+\mathrm{x}}$ ( B -filled $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type). The unit cell of B-filled $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type structure contains one octahedral column and two tetrahedral columns running infinitely along the short axis; the relative positioning of the columns implies that three equi-latitudinal trigonal prismatic voids per octahedron and tetrahedron are occupied by boron leaving unoccupied the voids adjacent to every next tetrahedron in the column. In both $\mathrm{Mn}_{7} \mathrm{C}_{3}{ }^{-}$and $\mathrm{Pd}_{6} \mathrm{CuB}_{3}$-type structures, the boron atoms fill two voids adjacent to one tetrahedron and one void adjacent to the consecutive tetrahedron along the column. While in the $\mathrm{Mn}_{7} \mathrm{C}_{3}$-type structure, one of every two octahedral columns is rotated around the short axis by $60^{\circ}$ with respect to the other, in $\mathrm{Pd}_{6} \mathrm{CuB}_{3}, 2 / 3$ of the octahedral columns are displaced giving rise to a new structure. The fourth member of the series of $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type related structures, the $\mathrm{Ca}_{7} \mathrm{Au}_{3}$ structure ${ }^{[12]}$ (and isotypic o- $\mathrm{Fe}_{7} \mathrm{C}_{3}$ compound recently discovered ${ }^{[11]}$ ) exhibits the arrangement of four octahedral columns in the unit cell, two of which are rotated around the short axis by $60^{\circ}$ with respect to the others. Thus the volumes of the HT $h-\left(\mathrm{Pt}_{0.717} \mathrm{Cu}_{0.283}\right)_{7} \mathrm{~B}_{3+\mathrm{x}}, o-\left(\mathrm{Pd}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathrm{~B}_{3}, h-\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}$ and insinuated from refinement of X -ray powder diffraction data the $\mathrm{Ca}_{7} \mathrm{Au}_{3}$-type derivative platinum copper boride relate as 1:2:3:4, respectively.

The increase of Pd content on progressing from the Cu -rich to Cu -poor members within the range of existence of $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type derivative structures $\left(\mathrm{Pd}_{58.5} \mathrm{Cu}_{11.5} \mathrm{~B}_{30^{-}}\right.$ $\mathrm{Pd}_{68.5} \mathrm{Cu}_{1.5} \mathrm{~B}_{30}$ at $650{ }^{\circ} \mathrm{C}$ ) is realized via continuous decrease of Cu occupancy at tetrahedral Cu 2 and Cu 5 sites ( $12 d, 6 c$ and $8 d, 4 c$ in $\mathrm{Pd}_{6} \mathrm{CuB}_{3}-$ and $\mathrm{Mn}_{7} \mathrm{C}_{3}$-type structures, respectively) up to complete vanishing of Cu in $4 c$ atom site of the orthorhombic unit cell of $\left(\mathrm{Pd}_{0.979} \mathrm{Cu}_{0.021}\right)_{7} \mathrm{~B}_{3}$. In the $\mathrm{Pt}-\mathrm{Cu}-\mathrm{B}$ system, the Pt-rich end (HT $h-\left(\mathrm{Pt}_{0.72-\mathrm{y}} \mathrm{Cu}_{0.28+\mathrm{y}}\right)_{7} \mathrm{~B}_{3+\mathrm{x}}, \mathrm{y} \cong 0-$ $0.07, \mathrm{x} \cong 0-0.3$ ) of the $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type derivative phase field in as cast conditions forms a small hexagonal lattice of B-filled $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type which on annealing at $600{ }^{\circ} \mathrm{C}$ decomposes to the low temperature form of $\mathrm{Pd}_{6} \mathrm{CuB}_{3}$-type. $\mathrm{Cu} / \mathrm{Pt}$ mixtures in the structures of both HT- and LTphases reside predominantly at all tetrahedral sites whereas the octahedral sites exhibit minor replacement of Pt by Cu . As inferred by Rietveld refinement of X-ray powder diffraction data and EDX analysis, the $\mathrm{Ca}_{7} \mathrm{Au}_{3}$-type derivative boride exists in the $\left(\mathrm{Pt}_{0.72-\mathrm{y}} \mathrm{Cu}_{0.28+\mathrm{y}}\right)_{7} \mathrm{~B}_{3+\mathrm{x}}$ ( $\mathrm{y} \cong 0.07-0.16, x \cong 0-0.3$ ) range both in as cast and annealed alloys.

Two new platinum containing boride systems $\mathrm{Pt}_{4.9} \mathrm{Cu}_{2.1} \mathrm{~B}_{3}\left(\mathrm{Pd}_{6} \mathrm{CuB}_{3}\right.$-type $)$ and $\mathrm{Pt}_{5.04} \mathrm{Cu}_{1.96} \mathrm{~B}_{3.3}$ (B-filled $\mathrm{Th}_{7} \mathrm{Fe}_{3}$-type) exhibit a superconducting transition at 0.67 K and 0.66 K , respectively. Such a small difference of the transition temperatures can not be accidental and most likely is a result of the influence of the common structural features of the two compounds. Considering this, one may treat the increase in the value of the upper critical field as the result of an increased amount of defects serving as pinning centers for the flux line lattice. Such a conclusion is supported by a 25 percent higher residual resistivity value found for $\mathrm{Pt}_{4.9} \mathrm{Cu}_{2.1} \mathrm{~B}_{3}$. However, Orlando et al. (1979) ${ }^{[43]}$ argue that the same mechanism should slightly decrease the superconducting transition temperature.

## Experimental Section

Synthesis, phase analysis and X-ray powder diffraction studies. Alloys were prepared from ingots of thoroughly re-melted pieces of palladium or platinum foil or sponge (Ögussa, Austria, 99.99 mass\%), Cu shot ( $2-6 \mathrm{~mm}, 99.999$ mass\%, ChemPur, Germany) and crystalline boron (ChemPur, Germany, 99.4 mass\%) by repeated arc melting under argon. The arcmelted buttons were wrapped in tantalum foil and vacuum-sealed in a quartz tube for annealing at $600^{\circ} \mathrm{C}\left(650{ }^{\circ} \mathrm{C}\right)$ for 240 hours. Lattice parameters and standard deviations were determined by least squares refinement of room temperature X-ray powder diffraction data obtained from a Guinier-Huber image plate employing monochromatic $\mathrm{CuK}_{\alpha 1}$ radiation and $G e$ as internal standard $\left(\mathrm{a}_{\mathrm{Ge}}=5.65791 \AA\right)$. Qualitative analysis has been performed from the results of Rietveld refinements of X-ray powder diffraction data (program FULLPROF ${ }^{[44]}$ ).

For precise quantitative analysis of the $\mathrm{Pd}(\mathrm{Pt}) / \mathrm{Cu}$ ratio, the samples were polished using standard procedure and were examined by scanning electron microscopy (SEM/EDX) using a Philips XL30 ESEM with EDAX XL-30 EDX-detector.

Electrical resistivity. Electrical resistivity was measured on $1 \times 1 \times 4 \mathrm{~mm}$ size bar specimen applying a standard 4-point technique with an a. c. bridge (Lakeshore 370 AC Resistance Bridge) in a temperature region from 0.3 to 300 K and in magnetic fields up to 1 T .

Single crystal X-ray diffraction studies. Single crystals suitable for X-ray diffraction studies were isolated from fragmented alloys. The crystals were measured on a four-circle Bruker APEX II diffractometer equipped with a CCD detector ( $\kappa$-geometry, Mo K $\alpha$ radiation); orientation matrices and unit cell parameters were derived with the APEX II software. ${ }^{[45]}$ Multi-scan absorption correction was applied using the program SADABS; frame data were reduced to intensity values applying the SAINT-Plus package. ${ }^{[46]}$ The structures were solved by direct methods and refined with the SHELXS-97 and SHELXL-97 programs, ${ }^{[47,48]}$ respectively. Further details concerning the experiments are summarized in Table 1.
Structure determination for $\mathbf{H T} \boldsymbol{h} \boldsymbol{-}\left(\mathbf{P t}_{0.717} \mathbf{C u}_{0.283}\right)_{7} \mathbf{B}_{3+\mathrm{x}}, \mathbf{x}=\mathbf{0 . 2 8}\left(\equiv \mathbf{P t}_{5.02} \mathbf{C u}_{1.98} \mathbf{B}_{3.28}\right)$. X-ray single crystal diffraction data indexing for the crystal selected from the as-cast alloy $\mathrm{Pd}_{45} \mathrm{Cu}_{20} \mathrm{~B}_{35}$ led to a hexagonal unit cell with lattice constants $\mathrm{a}=7.4424(12) \AA$ and $\mathrm{c}=4.8549(8) \AA$. The observed extinctions were compatible with the space groups $P 6_{3} / \mathrm{mmc}, P$ $62 c$ and $P 6_{3} m c$ (WinGX program package ${ }^{[49]}$ ) of which the noncentrosymmetric $P 6_{3} m c$ was confirmed to be correct by subsequent successful structure solution and refinement against $\mathrm{F}^{2}$. Structure solution applying direct methods resulted in three metal atom positions, one of which ( $6 c$ ) was assigned to Pt while the remaining two ( $6 c$ and $2 b$ ) showed considerably smaller electron densities; refinement of occupancy parameters, assuming mixed population for these atom sites, resulted in $\mathrm{M} 2=44.6 \% \mathrm{Pt} 2+55.4 \% \mathrm{Cu} 2$ and $\mathrm{M} 3=68.1 \% \mathrm{Pt} 3+31.9 \% \mathrm{Cu} 3$, respectively. Two boron atom sites $-6 c$ and $2 a$ - were located from the analysis of electron density peaks in difference Fourier maps according to reasonable interatomic distances between detected and proposed atoms, among them the $2 a$ site was found to be occupied for about $30 \%$. Refinement converged to a reliability factor as low as $\mathrm{R}_{\mathrm{F} 2}=0.0285$ applying anisotropic displacement parameters for the metal atoms and constrained isotropic displacement parameters for borons (Tables 1, 2, Table 1 Supporting Information).

Structure determination for $\boldsymbol{h}-\left(\mathbf{P d}_{\mathbf{0 . 8 6 4}} \mathbf{C u}_{\mathbf{0 . 1 3 6}}\right)_{7} \mathbf{B}_{\mathbf{3}}\left(\equiv \mathbf{P d}_{\mathbf{6 . 0 4 7}} \mathbf{C u}_{\mathbf{0 . 9 5 3}} \mathbf{B}_{3}\right)$. X-ray diffraction data for the crystal, which was selected from the alloy $\mathrm{Pd}_{58} \mathrm{Cu}_{10} \mathrm{~B}_{32}$ annealed at $600{ }^{\circ} \mathrm{C}$, were indexed with a primitive hexagonal unit cell with lattice parameters $\mathrm{a}=12.9426(9) \AA$ and $\mathrm{c}=4.8697(4) \AA$. Systematic absences in the diffraction data were consistent with three space
groups, $P 6_{3} / \mathrm{mcm}, P-6 c 2$ and $P 6_{3} c m$ (WinGX program package ${ }^{[49]}$ ), of which noncentrosymmetric $P 6_{3} \mathrm{~cm}$ proved to be correct during structure solution and refinement. Five palladium atom positions have been assigned and refined in a straightforward manner; two palladium sites exhibited relatively large displacement parameters suggesting that they are disordered with copper atoms ( $\mathrm{M} 2=0.627 \mathrm{Pd} 2+0.373 \mathrm{Cu} 2$ in $12 d$ and $\mathrm{M} 5=0.793 \mathrm{Pd} 5+0.207$ Cu 5 in $6 c$ ). Two boron sites were easily located in the difference Fourier map. Refinement of occupancy and displacement parameters proceeded successfully to a small residual value $\left(\mathrm{R}_{\mathrm{F} 2}=0.0263\right.$, max $/ \mathrm{min}$ residual electron density $1.59 /-1.51 \mathrm{el} / \mathrm{A}^{3}$ ) (Tables $1-2$, Table 1 Supporting Information).
Structure determination for $\boldsymbol{o}-\left(\mathbf{P d}_{0.93} \mathbf{C u}_{0.07}\right)_{7} \mathbf{B}_{3}\left(\equiv \mathbf{P d}_{6.504} \mathbf{C u}_{0.496} \mathbf{B}_{3}\right)$. The unit cell parameters of the single crystal obtained from the $\mathrm{Pd}_{65} \mathrm{Cu}_{5} \mathrm{~B}_{30}$ sample annealed at $600{ }^{\circ} \mathrm{C}$ (orthorhombic $P$, $\mathrm{a}=4.8971 \AA, \mathrm{~b}=7.5353 \AA, \mathrm{c}=12.9743 \AA$ ) and X-ray powder diffraction spectra recorded from both the annealed and as-cast alloys suggested isotypism with the $\mathrm{Mn}_{7} \mathrm{~B}_{3}$-type structure. ${ }^{[10]}$ Systematic extinctions in the single crystal X-ray data were consistent with two possible space group types: centrosymmetric Pnma (no. 62) and nonstandard setting of non-centrosymmetric $P n 2_{1} a$ (no. 33). The structure was solved in the centrosymmetric space group and direct methods provided 5 atom positions of palladium. The refinement of the structure with all sites fully occupied yielded rather large anisotropic displacement parameters for two Pd sites, Pd 2 and Pd 5 ; refinement of occupancy parameters assuming mixed population for these atom sites - resulted in $78.6 \% \mathrm{Pd} 2+21.4 \% \mathrm{Cu} 2$ and $93.2 \% \mathrm{Pd} 5+6.8 \% \mathrm{Cu} 5$, respectively. The atom positions of boron atoms were derived from difference Fourier synthesis. Refinement of the structure with anisotropic displacement parameters for metal atom sites converged to $\mathrm{R}=0.0247$ with residual electron densities smaller than $1.237 \mathrm{el} / \AA^{3}$ (Tables 1, 2, Table 1 Supporting Information).

Further details of the single crystal structure determination may be obtained from the Fachinformationszentrum Karlsruhe, 76344 Eggenstein Leopoldshafen, Germany (Fax: (+49)7247-808-666; e-mail: crysdata@fiz-karlsruhe.de), quoting the depository numbers CSD-.... (HT $h-\left(\mathrm{Pt}_{0.717} \mathrm{Cu}_{0.283}\right)_{7} \mathrm{~B}_{3+\mathrm{y}}(\mathrm{y}=0.28)$, CSD-.... $\left(h-\left(\mathrm{Pd}_{0.864} \mathrm{Cu}_{0.136}\right)_{7} \mathrm{~B}_{3}\right)$ and CSD-.... $(o-$ $\left.\left(\mathrm{Pd}_{0.93} \mathrm{Cu}_{0.07}\right)_{7} \mathrm{~B}_{3}\right)$.

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